

On the relations between principal eigenvalue and torsional rigidity

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“New Trends in PDE Constrained Optimization”
RICAM (Linz), October 14–18, 2019

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Our goal is to present some relations between two important quantities that arise in the study of elliptic equations. We always consider the **Laplace operator** $-\Delta$ with **Dirichlet** boundary conditions; other elliptic operator can be considered, while considering other boundary conditions (**Neumann** or **Robin**) adds to the problem severe extra difficulties, essentially due to the fact that in the Dirichlet case functions in $H_0^1(\Omega)$ can be easily extended to \mathbf{R}^d while this is not in general true in the other cases.

To better understand the two quantities we deal with, let us make the following two **measurements**.

- Take in Ω an uniform heat source ($f = 1$), fix an initial temperature $u_0(x)$, wait a long time, and measure the **average temperature** in Ω .
- Consider in Ω no heat source ($f = 0$), fix an initial temperature $u_0(x)$, and measure the decay rate to zero of the temperature in Ω .

The first quantity is usually called **torsional rigidity** and is defined as

$$T(\Omega) = \int_{\Omega} u \, dx$$

where u is the solution of

$$-\Delta u = 1 \text{ in } \Omega, \quad u \in H_0^1(\Omega).$$

In the thermal diffusion model $T(\Omega)/|\Omega|$ is the **average temperature** of a conducting medium Ω with **uniformly distributed heat sources** ($f = 1$).

The second quantity is the **first eigenvalue** of the Dirichlet Laplacian

$$\lambda(\Omega) = \min \left\{ \frac{\int_{\Omega} |\nabla u|^2 dx}{\int_{\Omega} u^2 dx} : u \in H_0^1(\Omega) \setminus \{0\} \right\}$$

In the thermal diffusion model, by the **Fourier analysis**,

$$u(t, x) = \sum_{k \geq 1} e^{-\lambda_k t} \langle u_0, u_k \rangle u_k(x),$$

so $\lambda(\Omega)$ represents the **decay rate** in time of the temperature when an initial temperature is given and no heat sources are present.

If we want, under the **measure constraint** $|\Omega| = m$, the highest average temperature, or the slowest decay rate, the optimal Ω is the same and is the ball of measure m . Also, **it seems consistent** to expect a slow (resp. fast) heat decay related to a high (resp. low) temperature. We then want to study if

$$\lambda(\Omega) \sim T^{-1}(\Omega),$$

or more generally

$$\lambda(\Omega) \sim T^{-q}(\Omega),$$

where by $A(\Omega) \sim B(\Omega)$ we mean

$$0 < c_1 \leq A(\Omega)/B(\Omega) \leq c_2 < +\infty \quad \text{for all } \Omega.$$

We further aim to study the so-called **Blasche-Santaló** diagram for the quantities $\lambda(\Omega)$ and $T(\Omega)$. This consists in identifying the set $E \subset \mathbf{R}^2$

$$E = \left\{ (x, y) : x = T(\Omega), y = \lambda(\Omega) \right\}$$

where Ω runs among the **admissible sets**. In this way, minimizing a quantity like

$$F(T(\Omega), \lambda(\Omega))$$

is reduced to the optimization problem in \mathbf{R}^2

$$\min \left\{ F(x, y) : (x, y) \in E \right\}.$$

The **difficulty** consists in the fact that characterizing the set E is **hard**. Here we only give some **bounds** by studying the inf and sup of

$$\lambda^\alpha(\Omega)T^\beta(\Omega)$$

when $|\Omega| = m$.

Since the two quantities scale as:

$$T(t\Omega) = t^{d+2}T(\Omega), \quad \lambda(t\Omega) = t^{-2}\lambda(\Omega)$$

it is **not restrictive** to reduce ourselves to the case $|\Omega| = 1$, which **simplifies** a lot the presentation.

For the relations between $T(\Omega)$ and $\lambda(\Omega)$:

- Kohler-Jobin ZAMP 1978;
- van den Berg, Buttazzo, Velichkov in Birkhäuser 2015
- van den Berg, Ferone, Nitsch, Trombetti Integral Equations Operator Theory 2016
- Lucardesi, Zucco paper in preparation;

The Blaschke-Santaló diagram has been studied for other pairs of quantities:

- for $\lambda_1(\Omega)$ and $\lambda_2(\Omega)$ by **D. Bucur, G. Buttazzo, I. Figueiredo** (SIAM J. Math. Anal. 1999);
- for $\lambda_1(\Omega)$ and $\text{Per}(\Omega)$ by **M. Dambrine, I. Ftouhi, A. Henrot, J. Lamboley** (paper in preparation).

For the inf/sup of

$$\lambda^\alpha(\Omega)T^\beta(\Omega)$$

the case $\beta = 0$ is well-known and reduces to the **Faber-Krahn** result (B ball with $|B| = 1$)

$$\min \{ \lambda(\Omega) : |\Omega| = 1 \} = \lambda(B),$$

while

$$\sup \{ \lambda(\Omega) : |\Omega| = 1 \} = +\infty$$

(take many **small balls** or a **long thin rectangle**).

Similarly, the case $\alpha = 0$ is also well-known through a **symmetrization argument** (**Saint-Venant** inequality):

$$\max \{T(\Omega) : |\Omega| = 1\} = T(B),$$

while

$$\inf \{T(\Omega) : |\Omega| = 1\} = 0$$

(take many **small balls** or a **long thin rectangle**).

The case when α and β have a different sign is also easy, since $T(\Omega)$ is increasing for the set inclusion, while $\lambda(\Omega)$ is decreasing.

So we can reduce the study to the case

$$\lambda(\Omega)T^q(\Omega)$$

with $q > 0$. If we want to remove the constraint $|\Omega| = 1$ the corresponding **scaling free** shape functional is

$$F_q(\Omega) = \frac{\lambda(\Omega)T^q(\Omega)}{|\Omega|^{(dq+2q-2)/d}}$$

that we consider on various classes of **admissible domains**.

We start by considering the class of **all domains** (with $|\Omega| = 1$). The known cases are:

- $q = 2/(d + 2)$ in which the minimum of $\lambda(\Omega)T^q(\Omega)$ is reached when Ω is a ball (**Kohler-Jobin** ZAMP 1978);
- $q = 1$ in which (**Pólya** inequality)

$$0 < \lambda(\Omega)T(\Omega) < 1.$$

When $0 < q \leq 2/(d+2)$:

$$\begin{cases} \min \lambda(\Omega)T^q(\Omega) = \lambda(B)T^q(B) \\ \sup \lambda(\Omega)T^q(\Omega) = +\infty. \end{cases}$$

For the minimum

$$\begin{aligned} \lambda(\Omega)T^q(\Omega) &= \lambda(\Omega)T(\Omega)^{2/(d+2)}T(\Omega)^{q-2/(d+2)} \\ &\geq \lambda(B)T(B)^{2/(d+2)}T(B)^{q-2/(d+2)} \\ &= \lambda(B)T^q(B), \end{aligned}$$

by Kohler-Jobin and Saint-Venant inequalities.

For the sup take $\Omega = N$ disjoint small balls.

When $2/(d+2) < q < 1$:

$$\begin{cases} \inf \lambda(\Omega)T^q(\Omega) = 0 \\ \sup \lambda(\Omega)T^q(\Omega) = +\infty. \end{cases}$$

For the sup take again $\Omega = N$ disjoint balls.

For the inf take as Ω the union of a fixed ball B_R and of N disjoint balls of radius ε .

We have

$$\lambda(\Omega)T^q(\Omega) = R^{-2}\lambda(B_1)T^q(B_1)\left(R^{d+2} + N\varepsilon^{d+2}\right)^q$$

and choosing first $\varepsilon \rightarrow 0$ and then $R \rightarrow 0$ we have that $\lambda(\Omega)T^q(\Omega)$ vanishes.

When $q = 1$:

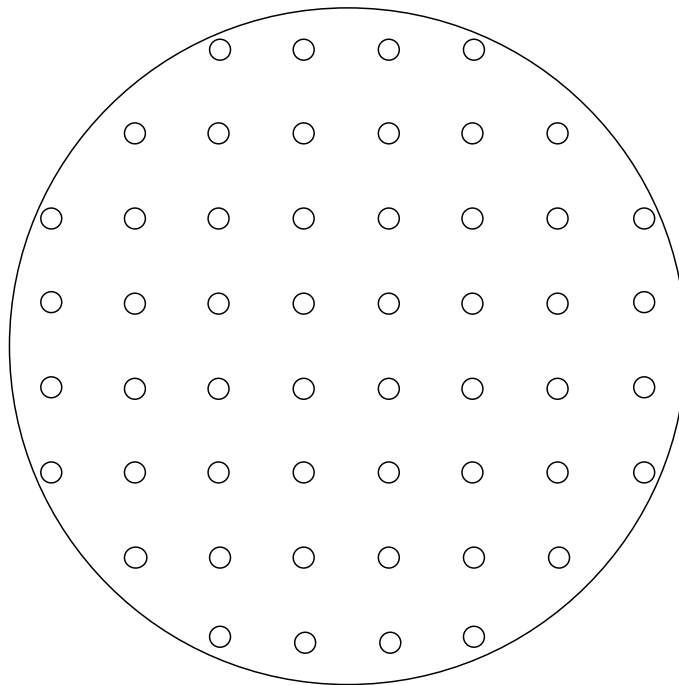
$$\inf \lambda(\Omega)T(\Omega) = 0, \quad \sup \lambda(\Omega)T(\Omega) = 1.$$

For the inf take as Ω the union of a fixed ball B_R and of N disjoint balls of radius ε , as above.

The sup equality, taking Ω a **finely perforated** domain, was shown by **van den Berg, Ferone, Nitsch, Trombetti** [Integral Equations Operator Theory 2016]. A shorter proof can be given using the theory of **capacitary measures**.

The finely perforated domains:

$\varepsilon =$ distance between holes $r_\varepsilon =$ radius of a hole
 $r_\varepsilon \sim \varepsilon^{d/(d-2)}$ if $d > 2$, $r_\varepsilon \sim e^{-1/\varepsilon^2}$ if $d = 2$.



When $q > 1$:

$$\inf \lambda(\Omega)T^q(\Omega) = 0, \quad \sup \lambda(\Omega)T^q(\Omega) < +\infty.$$

For the inf take as Ω the union of a fixed ball B_R and of N disjoint balls of radius ε , as above.

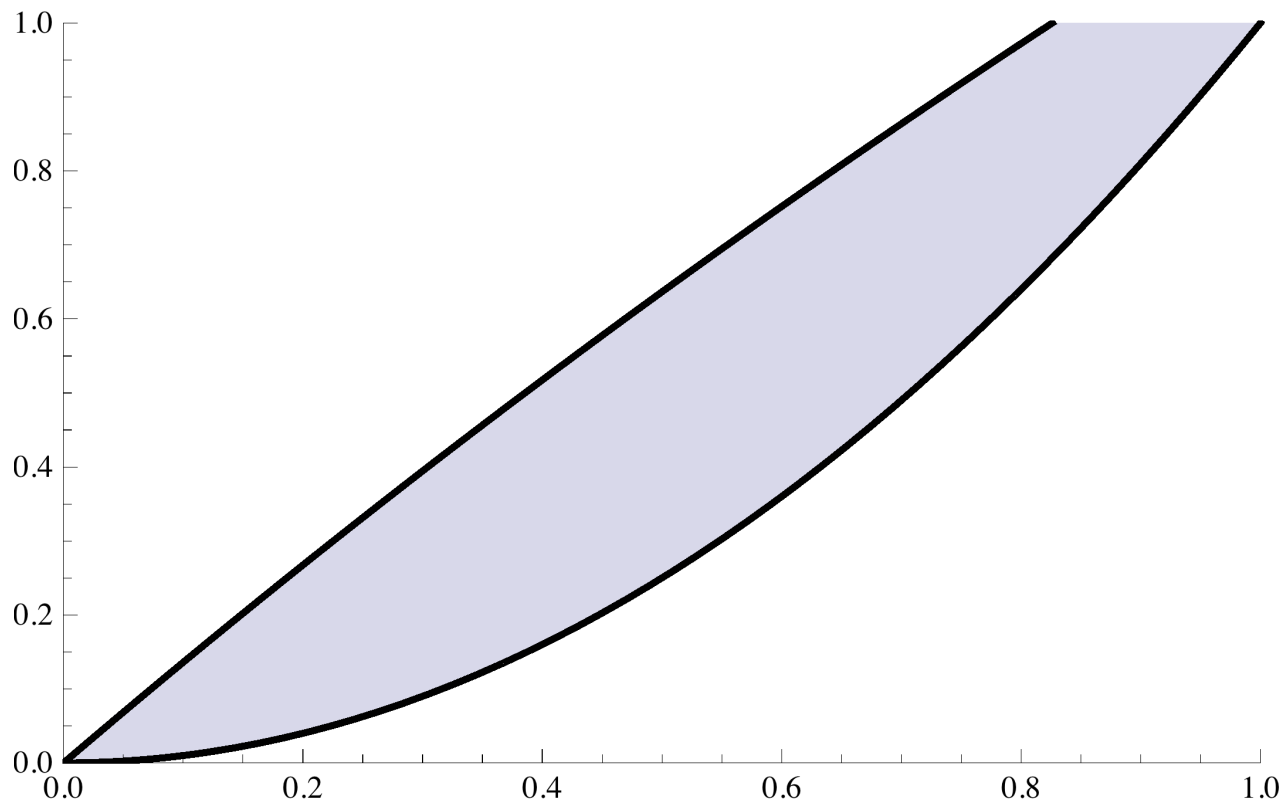
For the sup (using **Pólya** and **Saint-Venant**):

$$\begin{aligned} \lambda(\Omega)T^q(\Omega) &= \lambda(\Omega)T(\Omega)T^{q-1}(\Omega) \\ &\leq T^{q-1}(\Omega) \leq T^{q-1}(B) \end{aligned}$$

It would be interesting to **compute explicitly** $\sup F_q(\Omega)$ for $q > 1$ (**is it attained?**).

Summarizing: for **general domains** we have

	General domains Ω	
$0 < q \leq 2/(d + 2)$	$\min F_q(\Omega) = F_q(B)$	$\sup F_q(\Omega) = +\infty$
$2/(d + 2) < q < 1$	$\inf F_q(\Omega) = 0$	$\sup F_q(\Omega) = +\infty$
$q = 1$	$\inf F_q(\Omega) = 0$	$\sup F_q(\Omega) = 1$
$q > 1$	$\inf F_q(\Omega) = 0$	$\sup F_q(\Omega) < +\infty$



The Blaschke-Santaló diagram with $d = 2$, for $x = \lambda(B)/\lambda(\Omega)$ and $y = T(\Omega)/T(B)$ is **contained** in the colored region.

If we limit ourselves to consider only domains Ω that are union of disjoint disks of radii r_k we find

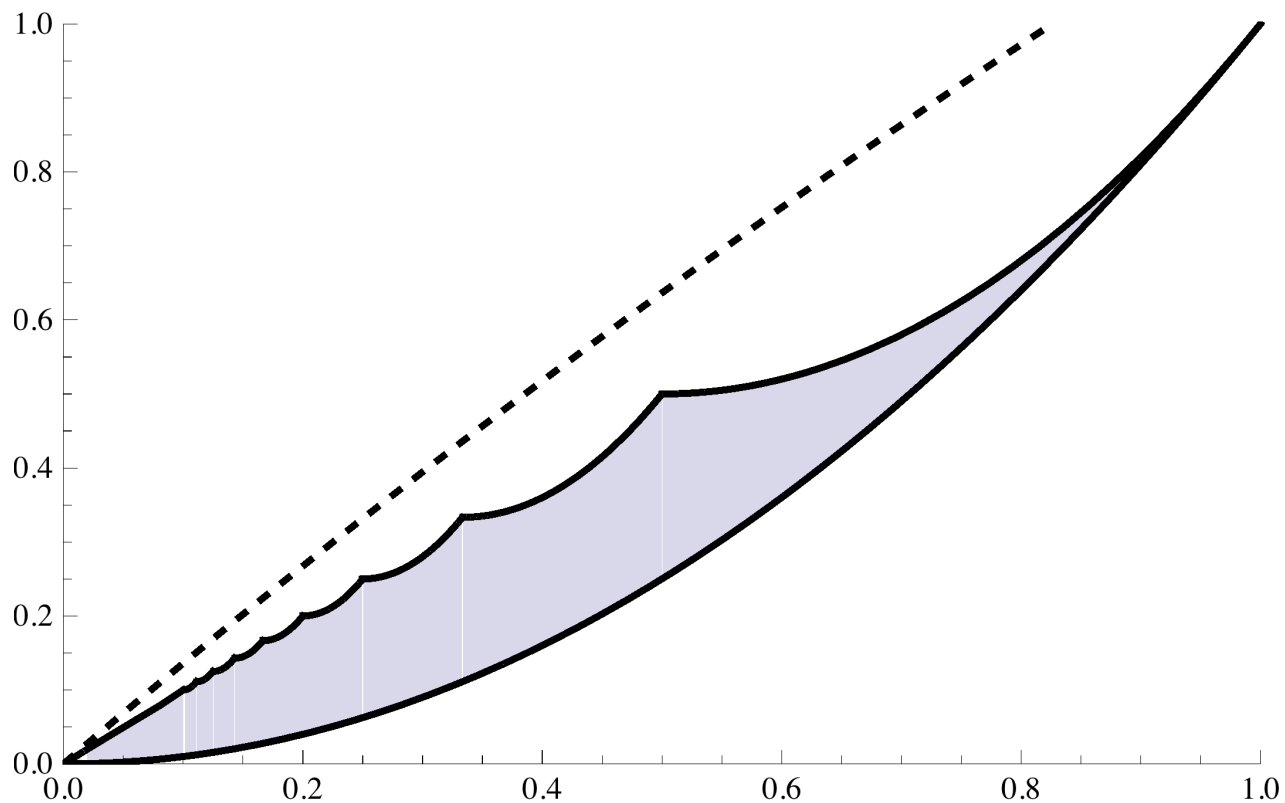
$$x = \frac{\max_k r_k^2}{\sum_k r_k^2}, \quad y = \frac{\sum_k r_k^4}{\left(\sum_k r_k^2\right)^2}.$$

It is not difficult to show that in this case we have

$$y \leq x^2[1/x] + \left(1 - x[1/x]\right)^2$$

where $[s]$ is the **integer part** of s .

In this way in the **Blaschke-Santaló** diagram we can reach the colored region in the picture below.



In the Blaschke-Santaló diagram with $d = 2$, the colored region can be reached by domains Ω made by union of disjoint disks.

The case $d = 1$

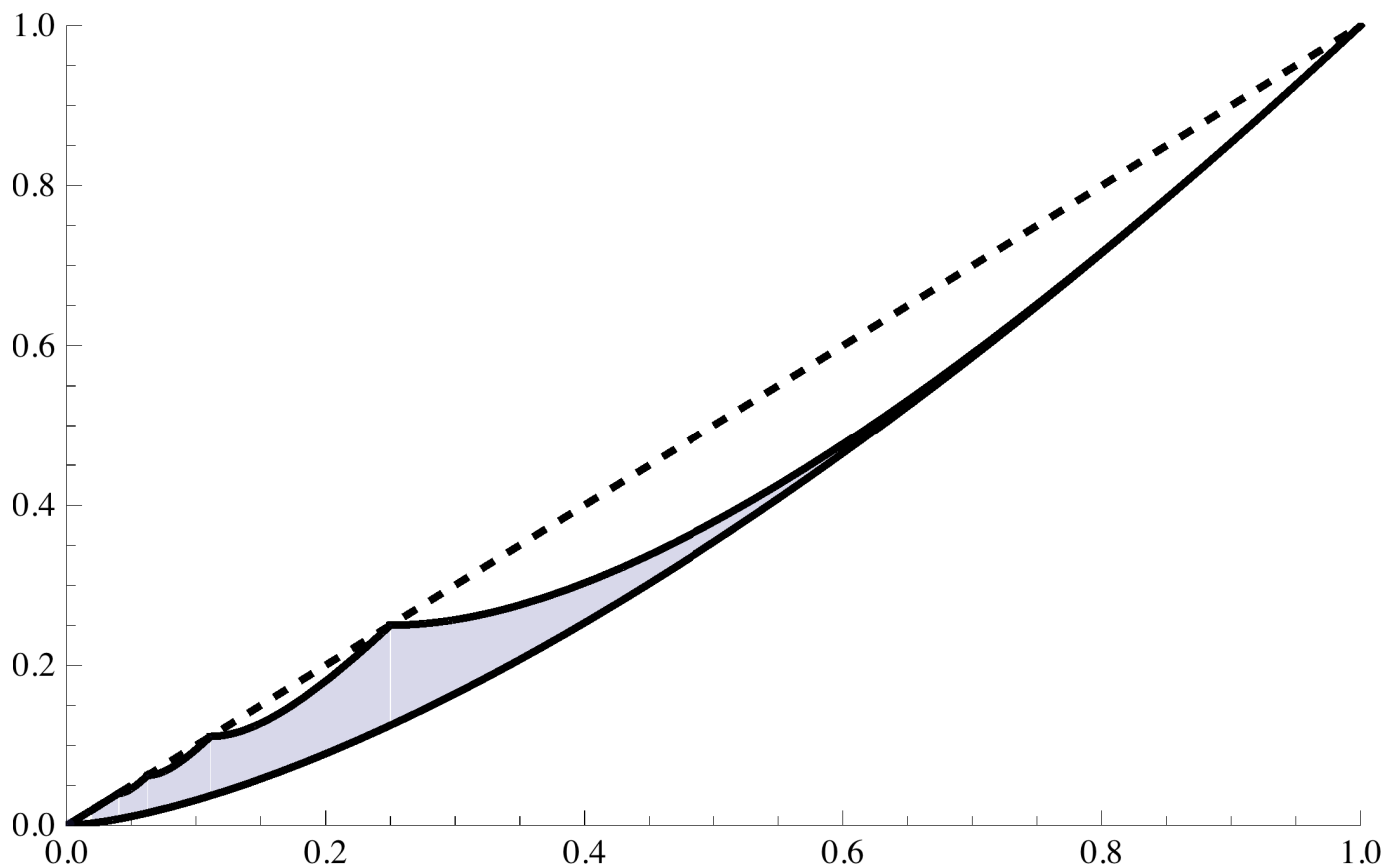
In the one-dimensional case every domain Ω is the union of disjoint intervals of half-length r_k , so that we have

$$x = \frac{\max_k r_k^2}{\left(\sum_k r_k\right)^2}, \quad y = \frac{\sum_k r_k^3}{\left(\sum_k r_k\right)^3}$$

and we deduce

$$y \leq x^{3/2} [x^{-1/2}] + \left(1 - x^{1/2} [x^{-1/2}]\right)^3$$

where $[s]$ is the **integer part** of s .



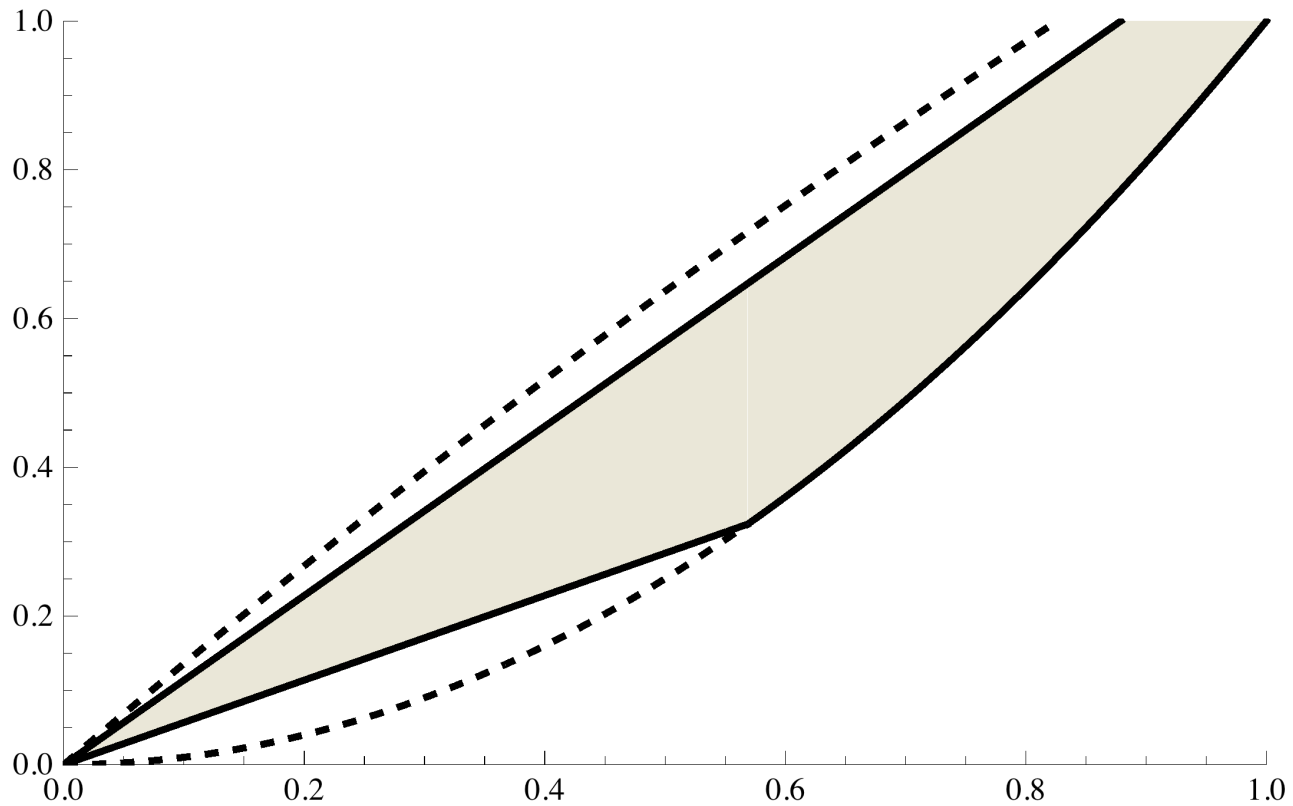
The **full** Blaschke-Santaló diagram in the case $d = 1$, where $x = \pi^2/\lambda(\Omega)$ and $y = 12T(\Omega)$.

The case Ω convex

If we consider only **convex** domains Ω , the Blaschke-Santaló diagram is clearly smaller. For the dimension $d = 2$ the **conjecture** is

$$\frac{\pi^2}{24} \leq \frac{\lambda(\Omega)T(\Omega)}{|\Omega|} \leq \frac{\pi^2}{12} \quad \text{for all } \Omega$$

where the left side corresponds to Ω a **thin triangle** and the right side to Ω a **thin rectangle**.



If the Conjecture for convex domains is true, the Blaschke-Santaló diagram is **contained** in the colored region.

At present the only available inequalities are the ones of [BFNT2016]: for every $\Omega \subset \mathbf{R}^2$ convex

$$0.2056 \approx \frac{\pi^2}{48} \leq \frac{\lambda(\Omega)T(\Omega)}{|\Omega|} \leq 0.9999$$

instead of the bounds provided by the conjecture, which are

$$\begin{cases} \pi^2/24 \approx 0.4112 & \text{from below} \\ \pi^2/12 \approx 0.8225 & \text{from above.} \end{cases}$$

In dimensions $d \geq 3$ the **conjecture** is

$$\frac{\pi^2}{2(d+1)(d+2)} \leq \frac{\lambda(\Omega)T(\Omega)}{|\Omega|} \leq \frac{\pi^2}{12}$$

- the right side asymptotically reached by a **thin slab**

$$\Omega_\varepsilon = \{(x', t) : 0 < t < \varepsilon\}$$

with $x' \in A_\varepsilon$, being A_ε a $d - 1$ dimensional ball of measure $1/\varepsilon$

- the left side asymptotically reached by a **thin cone** based on A_ε above and with height $d\varepsilon$.

Thin domains

We say that $\Omega_\varepsilon \subset \mathbf{R}^2$ is a **thin domain** if

$$\Omega_\varepsilon = \left\{ (s, t) : s \in]0, 1[, \varepsilon h_-(s) < t < \varepsilon h_+(s) \right\}$$

where ε is a small positive parameter and h_-, h_+ are two given (smooth) functions. We denote by $h(s)$ the **local thickness**

$$h(s) = h_+(s) - h_-(s)$$

and we assume that $h(s) \geq 0$.

The following **asymptotics** hold (as $\varepsilon \rightarrow 0$):

$$\lambda(\Omega_\varepsilon) \approx \frac{\varepsilon^{-2}\pi^2}{\|h\|_{L^\infty}^2} \quad [\text{Borisov-Freitas 2010}]$$

$$T(\Omega_\varepsilon) \approx \frac{\varepsilon^3}{12} \int h^3(s) ds$$

$$|\Omega_\varepsilon| \approx \varepsilon \int h(s) ds.$$

Hence, for a thin domain Ω_ε we have

$$\frac{\lambda(\Omega_\varepsilon)T(\Omega_\varepsilon)}{|\Omega_\varepsilon|} \approx \frac{\pi^2}{12} \frac{\int h^3(s) ds}{\|h\|_{L^\infty}^2 \int h ds}.$$

We are able to prove the [conjecture](#) above in the class of thin domains. More precisely:

- for every h we have

$$\frac{\int h^3(s) ds}{\|h\|_{L^\infty}^2 \int h ds} \leq 1 ;$$

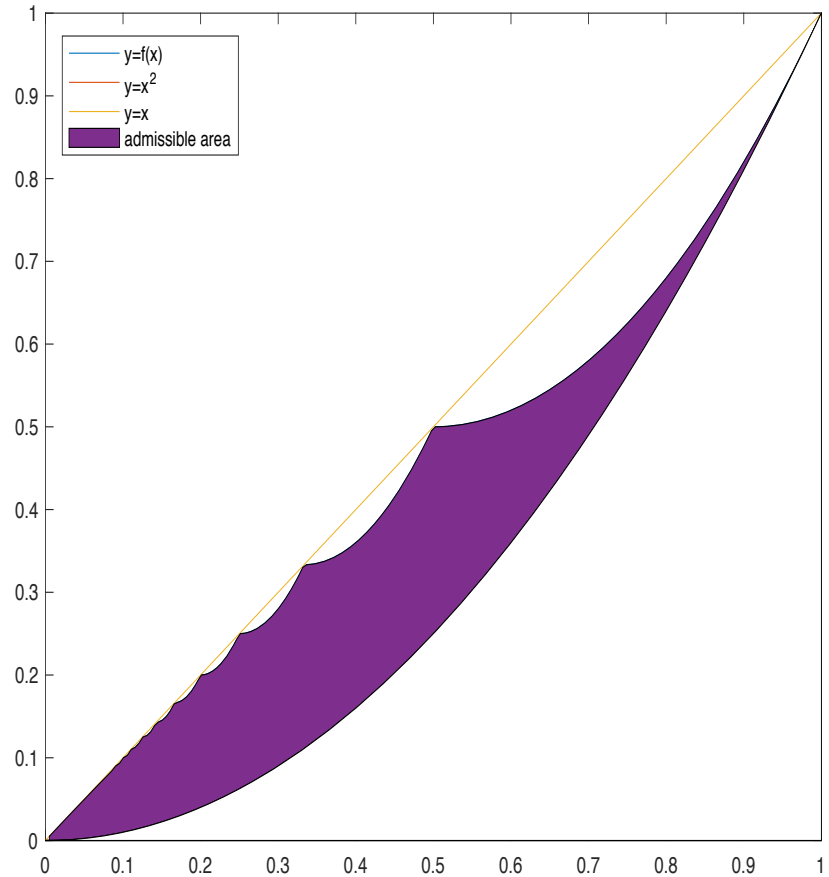
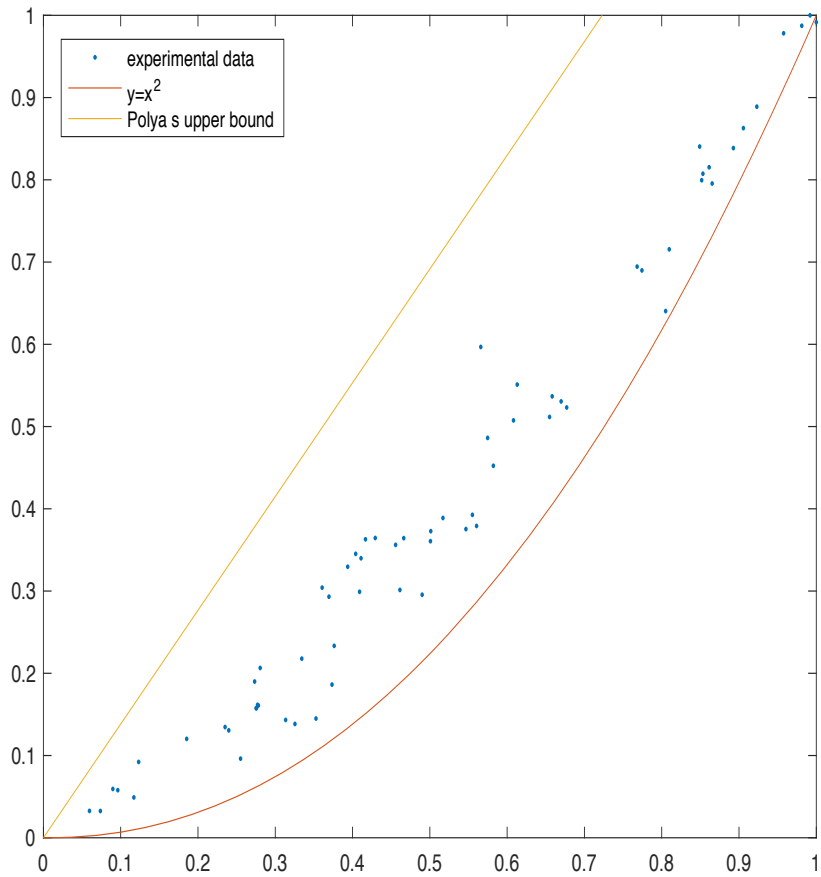
- for every h **concave** we have

$$\frac{\int h^3(s) ds}{\|h\|_{L^\infty}^2 \int h ds} \geq \frac{1}{2} .$$

Hence

$$\frac{\pi^2}{24} \leq \lim_{\varepsilon \rightarrow 0} \frac{\lambda(\Omega_\varepsilon) T(\Omega_\varepsilon)}{|\Omega_\varepsilon|} \leq \frac{\pi^2}{12}$$

where the right inequality holds for **all** thin domains, while the left inequality holds for **convex** thin domains.



Plot of 100 experimental domains (left), union of disks (right).

Open questions

- characterize $\sup_{|\Omega|=1} \lambda(\Omega)T^q(\Omega)$ when $q > 1$;
- prove (or disprove) the conjecture for convex sets;
- simply connected domains or star-shaped domains? The bounds may change;
- full Blaschke-Santaló diagram;
- p -Laplacian instead of Laplacian?
- efficient experiments (random domains?).